NUMERICAL METHODS IN LQG

aka: can we compute spin foam amplitudes?

NUMERICAL METHODS IN LQG

yes! let's see how.

DISCLAIMER



DISCLAIMER

- Very recent
- Requires a lot of time
- Solve "stinky" problems
- Lots of questions and very few answers



WHAT IS THIS LECTURE ABOUT?

- Who, What, Where? A little bit of how.
- 3D gravity. The Ponzano-Regge model.
- 4D gravity. The EPRL model.
- EPRL vertex amplitude computed with HPC methods



WHAT IS THIS LECTURE ABOUT?

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- EPRL vertex amplitude computed with HPC methods



On which topics(s) would you like lectures for the next LQG school? Perhaps a couple of lecture where show explicit calculations for transition amplitudes and such. I understand they are not exactly easy to do.

WHAT IS THIS LECTURE ABOUT?

I am sorry for my poor pictures



There will be a lot of happy little accidents

WHO, WHAT, WHERE?



- 3 years QG group Marseille
- divide and conquer! (\mathbf{P})
- HPC ready, fast and efficient
- X Truncation needed, error estimation, computationally expensive large triangulations
- $\int \frac{1}{2}$ single vertex, 3 vertices, bubble divergences
- 808 P.D., Francesco Gozzini, Giorgio Sarno + (Western U) Francesca Vidotto, Carlo Rovelli, Pietropaolo Frisoni, ...
 - Original paper <u>1807.03066</u>, advanced paper will appear soon applications: <u>1803.00835</u>, <u>1903.12624</u>, <u>2004.12911</u> and more to come

WHO, WHAT, WHERE?

Markov Chain Monte-Carlo Computation

- 1 year Florida Atlantic University + Fudan
- Integration over a Lefschetz Thimble to avoid oscillations
 - Huge spins are accessible
 - Not suitable for small spins
- ባር እስከ constraints and the second se
 - Main paper <u>2012.11515</u>

WHO, WHAT, WHERE?

Effective spin foam models

1 year - Bard College and Perimeter Institute

2 hours lecture yesterday

Seth Asante, Bianca Dittrich, Hal Haggard

Euclidean 2004.07013 and Lorentzian 2104.00485

Plus the hands-on experiment NEXT WEEK!

Renormalization and fixed points

6 years - U. Hamburg + U. Jena

- Fixed points under coarse graining.
- (m Truncation made of hypercubes (slightly generalized)

Benjamin Bahr and Sebastian Steinhaus

Original paper <u>1508.07961</u>, numerics <u>1605.07649</u>, review <u>2007.01315</u>

3D GRAVITY: THE PONZANO-REGGE MODEL (AGAIN)

Semiclassical limit of Racah coefficients - G.Ponzano and T.Regge

Transition amplitudes for Euclidean 3D LQG (discussed in length by Maïté, Edward Wilson-Ewing, and Valentin Bonzom)

3D GRAVITY: THE PONZANO-REGGE MODEL (AGAIN)

 $i_m^j = \left(\begin{array}{ccc} j_1 & j_2 & j_3 \\ m_1 & m_2 & m_3 \end{array}\right)$

Semiclassical limit of Racah coefficients - G.Ponzano and T.Regge

Transition amplitudes for Euclidean 3D LQG discussed in length by Maïté, Edward Wilson-Ewing, and Valentin Bonzom)

boundary state

Restrict to 3 valent spinnetworks (oriented graph, irreps, intertwiners)

Spins as lengths eigenvalues

Trivial intertwiners in 3D (1d space)

Nodes = Triangles (closure)

(Quantum) triangulations of a 2D surface (possibly curved)



Semiclassical limit of Racah coefficients - G.Ponzano and T.Regge

Transition amplitudes for Euclidean 3D LQG (discussed in length by Maïté and Edward Wilson-Ewing)

Choose a 2 complex dual to a 3D triangulation (vertices, edges, faces)



Semiclassical limit of Racah coefficients - G.Ponzano and T.Regge

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Semiclassical limit of Racah coefficients - G.Ponzano and T.Regge



3D GRAVITY: THE PONZANO-REGGE MODEL (WIRING)



SU(2) IN ONE SLIDE

Group matrix elements in the irrep j $\langle j,m| \ g \ |j,n\rangle \equiv D^j_{mn}(g)$

Numerics need a parametrization (Euler) - fundamental

 $g = e^{i\phi\frac{\sigma_3}{2}} e^{i\theta\frac{\sigma_3}{2}} e^{i\psi\frac{\sigma_3}{2}}$

Matrix elements explicit formula

$$D_{mn}^{j}(g) = e^{-i\phi \frac{m}{2}} e^{i\psi \frac{n}{2}} \sqrt{(j+m)!(j-m)!(j+n)!(j-n)!}$$

$$\sum_{s} (-1)^{m-n+s} \frac{\left(\cos \frac{\theta}{2}\right)^{2j+n-m-2s} \left(\sin \frac{\theta}{2}\right)^{m-n+2s}}{(j+n-s)!s!(m-n+s)!(j-m-s)!}$$

Integration of functions on the group using the Haar measure

 $\int f(g) \, dg = \frac{1}{16\pi^2} \int_0^{2\pi} d\phi \int_0^{\pi} d\theta \sin \theta \int_0^{4\pi} d\psi f(g)$

<u>A Primer of Group Theory for Loop Quantum Gravity and Spin-foams - Pierre Martin-Dussaud</u> <u>Quantum Theory of Angular Momentum - Varshalovich - aka the sacred text for SU(2) calculations</u>

Angular Momentum Algebra
$$\begin{split} & [J_a,J_b]=i\epsilon_{abc}J_c\\ & J_3\left|j,m\right\rangle=m\left|j,m\right\rangle\\ & J^2\left|j,m\right\rangle=j(j+1)\left|j,m\right\rangle \end{split}$$

Recipe to write any amplitude:

For each vertex

For each edge one group element g_e

For each face an irrep j_f and a matrix element $D_{m_sm_t}^{j_f}(g_s^{-1}g_t)$

Integrate over all the group elements $\int dg_e$

Sum over the internal faces spins with a weight $(2j_f + 1)$

Tell them about the interpretation in terms of reference systems



For each edge one group element g_e

For each face an irrep j_f and a matrix element $D_{m_sm_t}^{j_f}(g_s^{-1}g_t)$



 $D^{j_1}(g_d^{-1}g_c)D^{j_2}(g_d^{-1}g_b)D^{j_3}(g_d^{-1}g_a)$ $D^{j_4}(g_a^{-1}g_b)D^{j_5}(g_a^{-1}g_c)D^{j_6}(g_b^{-1}g_c)$

 $\int dg_e$

For each edge one group element g_e

For each face an irrep j_f and a matrix element $D_{m_sm_t}^{j_f}(g_s^{-1}g_t)$

Integrate over all the group elements

$$\int dg_a \, dg_b \, dg_c \, dg_d$$
$$D^{j_1}(g_d^{-1}g_c) D^{j_2}(g_d^{-1}g_b) D^{j_3}(g_d^{-1}g_a)$$
$$D^{j_4}(g_a^{-1}g_b) D^{j_5}(g_a^{-1}g_c) D^{j_6}(g_b^{-1}g_c)$$



 $\int dg_a \, dg_b \, dg_c \, dg_d$ $D^{j_1}(g_d^{-1}g_c) D^{j_2}(g_d^{-1}g_b) D^{j_3}(g_d^{-1}g_a)$ $D^{j_4}(g_a^{-1}g_b) D^{j_5}(g_a^{-1}g_c) D^{j_6}(g_b^{-1}g_c)$



Split the irreps, write the formula for the *D*, 6 sums of elementary integrals...

 $\int dg_a \, dg_b \, dg_c \, dg_d$ $D^{j_1}(g_d^{-1}g_c) D^{j_2}(g_d^{-1}g_b) D^{j_3}(g_d^{-1}g_a)$ $D^{j_4}(g_a^{-1}g_b) D^{j_5}(g_a^{-1}g_c) D^{j_6}(g_b^{-1}g_c)$

Or do the integrals analytically



$$\int dg D_{m_1 n_1}^{j_1}(g) D_{m_2 n_2}^{j_2}(g) D_{m_3 n_3}^{j_3}(g) = \begin{pmatrix} j_1 & j_2 & j_3 \\ m_1 & m_2 & m_3 \end{pmatrix} \begin{pmatrix} j_1 & j_2 & j_3 \\ n_1 & n_2 & n_3 \end{pmatrix} \longleftrightarrow = =$$

Each vertex gives you a Wigner $\{6 \not s\}$ ymbol.

Many explicit expressions: e.g. contraction of CG coefficients, in term of a sum of factorials, Hypergeometric function, etc.

Quantum Theory of Angular Momentum - Varshalovich - aka the sacred text for SU(2) calculations

11 A 12

Each vertex gives you a Wigner $\{6j\}$ ymbol.

$$\begin{cases} j_1 & j_2 & j_3 \\ j_4 & j_5 & j_6 \end{cases}$$

Many explicit expressions: e.g. contraction of CG coefficients, in term of a sum of factorials, Hypergeometric function, etc.

Quantum Theory of Angular Momentum - Varshalovich - aka the sacred text for SU(2) calculations

11 1 10

Numerically:

- 1. Code your own
- 2. Do not reinvent the wheel
 - In Mathematica is builtin SixJSymbol[{j1, j2, j3}, {j4, j5, j6}]
 - C, python and Matlab use <u>WIGXJPF</u>

Fast and Accurate Evaluation of Wigner 3j. 6j. and 9j Symbols Using Prime Factorization and Multiword Integer Arithmetic - H. T. Johansson and C. Forssén

3D GRAVITY: THE PONZANO-REGGE MODEL - THE ASYMPTOTIC

Semiclassical limit of Racah coefficients - G.Ponzano and T.Regge

Vertex amplitude for large spins is related to the Regge action $A_{v} = \begin{cases} \lambda j_{1} & \lambda j_{2} & \lambda j_{3} \\ \lambda j_{4} & \lambda j_{5} & \lambda j_{6} \end{cases} \approx \frac{1}{\lambda^{\frac{3}{2}}\sqrt{48\pi V}} \left(e^{i\lambda(\sum_{i} j_{i}\theta_{i}) + i\frac{\pi}{4}} + e^{-i\lambda(\sum_{i} j_{i}\theta_{i}) - i\frac{\pi}{4}} \right)$ Ji θ_i 0.15 See Hal Haggard and Seth Asante lecture for more details 0.10 **2** 0.05 0.00 Example: $j_i = 1$ Analogy with QM path -0.05integral 20 40 60 80 100 λ

3D GRAVITY: THE PONZANO-REGGE MODEL - THE ASYMPTOTIC

Vertex amplitude for large spins is related to the Regge action

$$A_{v} = \left\{ \begin{array}{c} \lambda j_{1} & \lambda j_{2} & \lambda j_{3} \\ \lambda j_{4} & \lambda j_{5} & \lambda j_{6} \end{array} \right\} \approx \frac{1}{\lambda^{\frac{3}{2}} \sqrt{48\pi V}} \left(e^{i\lambda(\sum_{i} j_{i}\theta_{i}) + i\frac{\pi}{4}} + e^{-i\lambda(\sum_{i} j_{i}\theta_{i}) - i\frac{\pi}{4}} \right)$$

$$\downarrow j_{i} \quad \theta_{i}$$

3D GRAVITY: THE PONZANO-REGGE MODEL - THE ASYMPTOTIC

Vertex amplitude for large spins is related to the Regge action

$$A_{v} = \begin{cases} \lambda j_{1} & \lambda j_{2} & \lambda j_{3} \\ \lambda j_{4} & \lambda j_{5} & \lambda j_{6} \end{cases} \approx \frac{1}{\lambda^{\frac{3}{2}}\sqrt{48\pi V}} \left(e^{i\lambda(\sum_{i} j_{i}\theta_{i})+i\frac{\pi}{4}} + e^{-i\lambda(\sum_{i} j_{i}\theta_{i})-i\frac{\pi}{4}} \right) \\ \downarrow j_{i} \\ \downarrow j_{i} \\ \theta_{i} \\ \downarrow j_{i} \\ \downarrow j_{i} \\ \theta_{i} \\ \psi_{i} \\ \psi_{i}$$









$$\sum_{x} (-1)^{x} (2x+1) \left\{ \begin{array}{cc} j_{5} & j_{8} & x \\ j_{9} & j_{6} & j_{1} \end{array} \right\} \left\{ \begin{array}{cc} j_{9} & j_{6} & x \\ j_{4} & j_{7} & j_{2} \end{array} \right\} \left\{ \begin{array}{cc} j_{4} & j_{7} & x \\ j_{8} & j_{5} & j_{3} \end{array} \right\}$$



$$P(X) = \sum_{x}^{X} (-1)^{x} (2x+1) \left\{ \begin{array}{cc} j_{5} & j_{8} & x \\ j_{9} & j_{6} & j_{1} \end{array} \right\} \left\{ \begin{array}{cc} j_{9} & j_{6} & x \\ j_{4} & j_{7} & j_{2} \end{array} \right\} \left\{ \begin{array}{cc} j_{4} & j_{7} & x \\ j_{8} & j_{5} & j_{3} \end{array} \right\}$$





3D GRAVITY: THE PONZANO-REGGE MODEL - THE BUBBLE



LET'S HAVE A SHORT BREAK

<u>LQG vertex with finite Immirzi parameter - Jonathan Engle, Etera Livine, Roberto Pereira, Carlo Rovelli</u> <u>Covariant Loop Quantum Gravity - Carlo Rovelli, Francesca Vidotto</u>

Transition amplitudes for Lorentzian 4D LQG (why is defined like this? - Maïté)

<u>LQG vertex with finite Immirzi parameter - Jonathan Engle, Etera Livine, Roberto Pereira, Carlo Rovelli</u> <u>Covariant Loop Quantum Gravity - Carlo Rovelli, Francesca Vidotto</u>

Transition amplitudes for Lorentzian 4D LQG why is defined like this? - Maïté)

boundary state

Restrict to 4 valent spinnetworks (KKL extension) Spins as areas eigenvalues Intertwiners in 4D = Quantum tetrahedra (closure) (Quantum) triangulations of a 3D surface (twisted geometries) i_4 $i = \left(\begin{array}{c} j_1 & j_2 & j_3 & j_4 \\ m_1 & m_2 & m_3 & m_4 \end{array} \right)^{(j_{12})}$



4 VALENT INTERTWINER = QUANTUM TETRAHEDRON

$$i = \begin{pmatrix} j_1 & j_2 & j_3 & j_4 \\ m_1 & m_2 & m_3 & m_4 \end{pmatrix}^{(j_{12})} = \sum_m (-1)^{j_{12}-m} \begin{pmatrix} j_1 & j_2 & j_{12} \\ m_1 & m_2 & m \end{pmatrix} \begin{pmatrix} j_{12} & j_3 & j_4 \\ -m & m_3 & m_4 \end{pmatrix}$$

Areas:

$$\left|i\right\rangle = j_i(j_i+1)\left|i\right\rangle$$

 J_i^2

Closure:
$$|i\rangle \in \operatorname{Inv}\left(H^{(j_1)} \otimes H^{(j_2)} \otimes H^{(j_3)} \otimes H^{(j_4)}\right) \rightarrow \left(\vec{J_1} + \vec{J_2} + \vec{J_3} + \vec{J_4}\right)|i\rangle = 0$$

 $\vec{J}_1 \cdot \vec{J}_2 |i\rangle = \frac{j_{12}(j_{12}+1) - j_1(j_1+1) - j_2(j_2+1)}{2} |i\rangle$

D. angle:

0.00004

0.00003

0.00002

0.00001

 $|\langle j_{12}|c(\vec{n}_i)\rangle|^2$



 $|\langle j_{13}|c(\vec{n}_i)\rangle|^2$



Coherent intertwiner:

 $\underbrace{j_{1}}_{15} \quad c_{i}(\vec{n_{i}}) = \begin{pmatrix} j_{1} & j_{2} & j_{3} & j_{4} \\ m_{1} & m_{2} & m_{3} & m_{4} \end{pmatrix}^{(j_{12})} \langle j_{1}m_{1}|j_{1}\vec{n}_{1}\rangle \langle j_{2}m_{2}|j_{2}\vec{n}_{2}\rangle \langle j_{3}m_{3}|j_{3}\vec{n}_{3}\rangle \langle j_{4}m_{4}|j_{4}\vec{n}_{4}\rangle$



<u>LQG vertex with finite Immirzi parameter - Engle, Livine, Pereira, Rovelli</u> <u>Covariant Loop Quantum Gravity - Rovelli, Vidotto</u> Maïté's lectures at this school

Recipe to write any amplitude:

Choose a 2 complex dual to a 4D triangulation (vertices, edges, faces)





SL(2,C) IN ONE SLIDE

Group matrix elements in the irrep
$$(\rho, k)$$
 $\langle \rho, k ; l, n | h | \rho, k ; j, m \rangle = D_{lnjm}^{\rho,k}(h)$
 $l, j \ge k$

Numerics need a parametrization - many - Cartan $h = ue^{\frac{r}{2}\sigma_3}v^{\dagger}$ $u, v \in SU(2), r \in \mathbb{R}^+$

Integration of functions on the group using the Haar measure

$$\int dh = \int du dv \sinh(r) dr$$

$$\begin{aligned} \text{Matrix elements explicit formula} \quad d_{jlm}^{(\rho,k)}(r) &= (-1)^{j-l} \sqrt{\frac{(i\rho - j - 1)! \, (j + i\rho)!}{(i\rho - l - 1)! \, (l + i\rho)!}} \frac{\sqrt{(2j + 1)(2l + 1)}}{(j + l + 1)!} e^{(i\rho - k - m - 1)r} \\ &= \sum_{o} D_{no}^{l}(u) d_{ljo}^{\rho,k}(r) D_{om}^{j}(v^{\dagger}) \\ &= \sum_{s,t} (-1)^{s+t} e^{-2tr} \frac{(k + s + m + t)! (j - k)! (l - m)!}{t! s! (j - k - s)! (j - m - s)! (k + m + s)! (l - k - t)! (l - m - t)! (k + m + t)!} \\ &= \sum_{o} D_{no}^{l}(u) d_{ljo}^{\rho,k}(r) D_{om}^{j}(v^{\dagger}) \\ &= \sum_{s,t} (-1)^{s+t} e^{-2tr} \frac{(k + s + m + t)! (j + l - k - m - s - t)!}{t! s! (j - k - s)! (j - m - s)! (k + m + s)! (l - k - t)! (l - m - t)! (k + m + t)!} \\ &= 2F_1 \left[\{l - i\rho + 1, k + m + s + t + 1\}, \{j + l + 2\}; 1 - e^{-2r} \right] \end{aligned}$$

<u>A Primer of Group Theory for Loop Quantum Gravity and Spin-foams - Pierre Martin-Dussaud</u> <u>The Lorentz group and harmonic analysis - W. Rühl</u>

$$\begin{split} \textbf{SL(2,C) Algebra} \\ \hline & [L_i, L_j] = i\epsilon_{ijk}L_k \\ & [L_i, K_j] = i\epsilon_{ijk}K_k \\ & [K_i, K_j] = -i\epsilon_{ijk}L_k \\ \end{split} \\ \textbf{Unit. Irreps: } \rho \in \mathbb{R} \quad k \in \mathbb{Z}/2 \\ (K^2 - L^2) |\rho, k; \ j, m\rangle = (\rho^2 - k^2 + 1) |\rho, k; \ j, m\rangle \\ & \vec{K} \cdot \vec{L} |\rho, k; \ j, m\rangle = \rho k |\rho, k; \ j, m\rangle \\ & L^2 |\rho, k; \ j, m\rangle = j(j+1) |\rho, k; \ j, m\rangle \\ & L_3 |\rho, k; \ j, m\rangle = m |\rho, k; \ j, m\rangle \end{split}$$

<u>LQG vertex with finite Immirzi parameter - Engle, Livine, Pereira, Rovelli</u> <u>Covariant Loop Quantum Gravity - Rovelli, Vidotto</u> Maïté's lectures at this school

Impose simplicity constraints as a restriction on irreps

 $|j,m
angle \xrightarrow{Y_{\gamma}} |\gamma j,j; j,m
angle$

 γ Immirzi parameter

Impose simplicity constraints as a restriction on irreps

 $|j,m
angle \xrightarrow{Y_{\gamma}} |\gamma j,j\,;\,\,j,m
angle$

Recipe to write any amplitude:

For each vertex:

For each edge one group element h_e

For each face an SU(2) spin j_f and a matrix element $D_{j_f n j_f m}^{\gamma j_f, j_f}(h_s^{-1}h_t)$

$$D_{j_f n \, j_f m}^{\gamma j_f, \, j_f}(h_s^{-1} h_t) = \sum_{l=j_f}^{\infty} \sum_{o=-l}^{l} D_{j_f n \, lo}^{\gamma j_f, \, j_f}(h_s^{-1}) D_{lo \, j_f m}^{\gamma j_f, \, j_f}(h_t)$$

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 γ Immirzi parameter

Perform the integrals over SL(2,C) using the Cartan decomposition:



<u>LQG vertex with finite Immirzi parameter - Engle, Livine, Pereira, Rovelli</u> <u>Covariant Loop Quantum Gravity - Rovelli, Vidotto</u> Maïté's lectures at this school

 (k_e)

$$A_{v}(j_{f}, i_{e}) = \sum_{l_{f}, k_{e}} \left(\prod_{e} d_{k_{e}} B_{4}(j_{f}, l_{f}, i_{e}, k_{e}) \right) \{15j\}(l_{f}, k_{e})$$
$$B_{4}(l_{f}, j_{f}, i_{e}, k_{e}) \equiv \sum_{m_{f}} {\binom{l_{f}}{m_{f}}}^{(i_{e})} \left(\int_{0}^{\infty} d\mu(r) \prod_{f} d_{l_{f}j_{f}m_{f}}^{(\gamma j_{f}, j_{f})}(r) \right) {\binom{j_{f}}{m_{f}}}$$



SL2CFOAM



- Written in C
- SU(2) invariants WIGXJPF
- Arbitrary precision arithmetic
- parallelizable
 - First attempt, good as long as it works
- largely unoptimized, slow

SL2CFOAM - NEXT



<u>sl2cfoam-next</u> - Francesco Gozzini (the library is available on git, the paper will appear soon)

- Written in C
- SU(2) invariants <u>WIGXJPF</u>
- Arbitrary precision arithmetic
- Optimized to be efficient
 - (boosters, data structures)
- HPC ready (OMP, MPI, GPU)
- Julia interface

SL2CFOAM - NEXT



<u>sl2cfoam-next</u> - Francesco Gozzini (the library is available on git, the paper will appear soon)

User friendly!

using SL2Cfoam

init SL2Cfoam library Immirzi = 0.1; folder = "/home/francesco/phd/src/data_sl2cfoam_next"; conf = SL2Cfoam.config(VerbosityOff, NormalAccuracy, 100, 0); SL2Cfoam.cinit(folder, Immirzi, conf);

compute vertex tensor for increasing shells $(j = 1, \Delta s = 0 \rightarrow 4)$ vs = [vertex_compute(ones(10), s) for s in 0:4];

print values value with intertwiners (0,0,0,0,0)
[v.a[1, 1, 1, 1, 1] for v in vs]

5-element Array{Float64,1}: 5.992940521928846e-8 8.77026353463005e-8 1.0511153396261617e-7 1.161678351516197e-7 1.2344027926668653e-7

harder better faster



pics stolen from Francesco's ILQGS seminar

EPRL VERTEX AMPLITUDE : DIVIDE AND CONQUER

$$A_{v}(j_{f}, i_{e}) = \sum_{l_{f}, k_{e}} \left(\prod_{e} d_{k_{e}} B_{4}(j_{f}, l_{f}, i_{e}, k_{e}) \right) \{15j\}(l_{f}, k_{e})$$
$$B_{4}(l_{f}, j_{f}, i_{e}, k_{e}) \equiv \sum_{m_{f}} {\binom{l_{f}}{m_{f}}}^{(i_{e})} \left(\int_{0}^{\infty} d\mu(r) \prod_{f} d_{l_{f}j_{f}m_{f}}^{(\gamma j_{f}, j_{f})}(r) \right) {\binom{j_{f}}{m_{f}}}^{(i_{e})}$$

This is the simple part

Fast and Accurate Evaluation of Wigner 3j. 6j. and 9j Symbols Using Prime Factorization and Multiword Integer Arithmetic - H. T. Johansson and C. Forssén



EPRL VERTEX AMPLITUDE : DIVIDE AND CONQUER

$$A_{v}(j_{f}, i_{e}) = \sum_{l_{f}, k_{e}} \left(\prod_{e} d_{k_{e}} B_{4}(j_{f}, l_{f}, i_{e}, k_{e}) \right) \{15j\}(l_{f}, k_{e})$$
$$B_{4}(l_{f}, j_{f}, i_{e}, k_{e}) \equiv \sum_{m_{f}} {\binom{l_{f}}{m_{f}}}^{(i_{e})} \left(\int_{0}^{\infty} d\mu(r) \prod_{f} d_{l_{f}j_{f}m_{f}}^{(\gamma j_{f}, j_{f})}(r) \right) {\binom{j_{f}}{m_{f}}}$$

Relatively complicated



 (k_e)

THE BOOSTER FUNCTIONS

Encode the details of the model (imposition of the simplicity constraints) Well understood semiclassical limit <u>Asymptotics of SL(2,C) coherent invariant</u>

Asymptotics of SL(2,C) coherent invariant tensors - P.D., Fanizza, Martin-Dussaud, Speziale

$$B_4(l_f, j_f, i_e, k_e) \equiv \sum_{m_f} \binom{l_f}{m_f}^{(i_e)} \left(\int_0^\infty d\mu(r) \prod_f d_{l_f j_f m_f}^{(\gamma j_f, j_f)}(r) \right) \binom{j_f}{m_f}^{(k_e)}$$

We know the analytic result of the integral. Related to SL(2,C) Clebsch-Gordan coefficient.

Boosting Wigner's nj-symbols - Simone Speziale

Not efficient to implement numerically.

THE BOOSTER FUNCTIONS

Encode the details of the model (imposition of the simplicity constraints) Well understood semiclassical limit <u>Asymptotics of SL(2,C) coherent invariant</u>

Asymptotics of SL(2,C) coherent invariant tensors - P.D., Fanizza, Martin-Dussaud, Speziale

$$B_4(l_f, j_f, i_e, k_e) \equiv \sum_{m_f} \binom{l_f}{m_f}^{(i_e)} \left(\int_0^\infty d\mu(r) \prod_f d_{l_f j_f m_f}^{(\gamma j_f, j_f)}(r) \right) \binom{j_f}{m_f}^{(k_e)}$$

Unbounded? Not really!

$$\begin{aligned} d_{jlm}^{(\rho,k)}(r) &= (-1)^{j-l} \sqrt{\frac{(i\rho - j - 1)! (j + i\rho)!}{(i\rho - l - 1)! (l + i\rho)!}} \frac{\sqrt{(2j + 1)(2l + 1)}}{(j + l + 1)!} e^{(i\rho - k - m - 1)r} \\ &\sqrt{(j + k)! (j - k)! (j + m)! (j - m)! (l + k)! (l - k)! (l - m)!} \\ &\sum_{s,t} (-1)^{s+t} e^{-2tr} \frac{(k + s + m + t)! (j + l - k - m - s - t)!}{l! s! (j - k - s)! (j - m - s)! (k + m + s)! (l - k - t)! (l - m - t)! (k + m + t)!} \\ &_2F_1 \left[\{l - i\rho + 1, k + m + s + t + 1\}, \{j + l + 2\}; 1 - e^{-2r} \right] \end{aligned}$$

THE BOOSTER FUNCTIONS

Encode the details of the model (imposition of the simplicity constraints) Well understood semiclassical limit <u>Asymptotics of SL(2,C) coherent invarian</u>

Asymptotics of SL(2,C) coherent invariant tensors - P.D., Fanizza, Martin-Dussaud, Speziale

$$B_4(l_f, j_f, i_e, k_e) \equiv \sum_{m_f} \binom{l_f}{m_f}^{(i_e)} \left(\int_0^\infty d\mu(r) \prod_f d_{l_f j_f m_f}^{(\gamma j_f, j_f)}(r) \right) \binom{j_f}{m_f}^{(k_e)}$$

High frequency oscillations!



The result is in cancellations!

- Arbitrary precision routines (long double has not enough precision)
- Adaptive Gauss-Kronrod quadrature method for numerical integration (1D, sampling not known in advance, recycle points) - adapted from GSL source

THE SUMS

sl2cfoam-next - Francesco Gozzini (to appear soon)

$$A_{v}(j_{f}, i_{e}) = \sum_{l_{f}, k_{e}} \left(\prod_{e} d_{k_{e}} B_{4}(j_{f}, l_{f}, i_{e}, k_{e}) \right) \{15j\}(l_{f}, k_{e})$$

Store everything in multidimensional arrays

Sums thought as tensors contractions. Optimized routines (BLAS) instead of caveman loops

Same technique to contract vertices

THE SUMS - THE DOWNSIDE

$$A_v(j_f, i_e) = \sum_{l_f, k_e}^{O} \left(\prod_e d_{k_e} B_4(j_f, l_f, i_e, k_e) \right) \{15j\}(l_f, k_e)$$

$$D_{j_f n j_f m}^{\gamma j_f, j_f}(h_s^{-1} h_t) = \sum_{l=j_f}^{\infty} \sum_{o=-l}^{l} D_{j_f n lo}^{\gamma j_f, j_f}(h_s^{-1}) D_{lo j_f m}^{\gamma j_f, j_f}(h_t)$$

To perform numerical computation we need to cutoff these sums

$$\sum_{l=j_f}^{\infty} \approx \sum_{l=j_f}^{j_f + \Delta s}$$

THE SUMS - THE DOWNSIDE $A_v(j_f, i_e) = \sum_{l_f=j_f}^{j_f+\Delta s} \sum_{k_e} \left(\prod_e d_{k_e} B_4(j_f, l_f, i_e, k_e)\right) \{15j\}(l_f, k_e)$

6 sums. We decided to impose uniformly the cutoff

- We know the sums are convergent
- Slang: Δs shells
- Good approximation (eventually)
- Empirical error of the truncation (bad)
- depends on everything



Pietropaolo Frisoni, Carlo Rovelli, Francesco Gozzini, P. D.

Strategy: as much as possible! Error?



Similar to P.R. EPRL amplitude with bdr coherent states peaked on the bdr of a Lorentzian 4-simplex

- power law behavior
- oscillation frequency equal to Regge Action of the Lorentzian 4-simplex



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APPLICATIONS - 3 VERTEX AMPLITUDE - 1 INTERNAL FACE



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APPLICATIONS - IR (LARGE-VOLUME) DIVERGENCES

Two results in the literature, IR melonic divergences, necessary step for renormalization

- bubbles are divergent
- **bounded from below** $\log K$
- bounded from above K^9
- numerical estimation is in progress

(the same kind in Valentin Bonzom's lecture)

<u>Self-Energy of the Lorentzian EPRL-FK Spin Foam Model of Quantum Gravity - Aldo Riello</u> <u>Infrared divergences in the EPRL-FK Spin Foam model - P.D.</u>

w.i.p. Francesco Gozzini, Pietropaolo Frisoni, Carlo Rovelli, Francesca Vidotto



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CONCLUSION

Doing computations with the Ponzano Regge model is an excellent way to play with spin foams on your laptop. (more minor technical difficulties of 4D, shares formal aspects, extension is not straightforward).

We have a well-written, optimized, parallelizable, user-friendly code (sl2cfoam-next) that allows computing EPRL Lorentzian spin foam amplitudes. (Elementary ones on your laptop, to do serious computations, HPC is needed - our resources are orders of magnitudes smaller than LQCD guys ones)

Field in constant development. We change idea very quickly! We should not be afraid to.

If you want to contribute you are very welcome to fork the repository and have fun!

- Implement Euclidean EPRL model (bounded sums)
- Extrapolation in the internal spin summation (truncation control) (Pietropaolo Frisoni, Francesco Gozzini)
- Amplitudes representing physical processes

(Carlo Rovelli, Farshid Soltani and Francesca Vidotto @Western U. on bouncing black holes)

(P.D., Francesco Gozzini, Alessandro Nicotra)