Quantum info of LQG states



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Collaboration with Eugenio Bianchi Phys. Rev. D 100, 105010 (2019) Phys. Rev. D 99, 086013 (2019) and w.i.p. Plan of the talk:

Motivations and background

Typical entropy: Page curve and its variance

Entropy and correlations in Bell-network states

Correlations at space-like separation

Quantum field theory \longrightarrow Fock space contains:

(i) states with no space-like correlations

(ii) states with specific short-ranged correlations





(e.g. Minkowski vacuum)

 $\langle \phi(0)\phi(r)\rangle\approx r^{-2}$

Loop quantum gravity (i) states with no space-like correlations (spin-networks) (ii) states with specific short-ranged correlations (many proposals)

LQG Hilbert space

 $\mathcal{H}_{\Gamma} = \bigoplus_{j_{\ell}} \bigotimes_n \mathcal{H}_n$



LQG Hilbert space





 $\mathcal{H}_n = \mathrm{Inv}\left(j_1 \otimes \cdots \otimes j_F\right)$

Quantum Polyhedra

[Bianchi, P.D., Speziale PRD 2010]





Geometric picture from LQG states

Information-theoretic bounds on correlations

State

 $|\psi\rangle \in \mathcal{H}_{\Gamma}$

Subsystem A

 $O_A \in \mathcal{A}_A$

Correlations

 $\mathcal{C}(O_A, O_B) = \langle \psi | O_A O_B | \psi \rangle - \langle \psi | O_A | \psi \rangle \langle \psi | O_B | \psi \rangle$

Bounded by mutual information

$$\frac{1}{2} \left(\frac{\mathcal{C}\left(O_A, O_B\right)}{\|O_A\| \|O_B\|} \right)^2 \leq S_A(\psi) + S_B(\psi) - S_{AB}(\psi) - S_{AB}(\psi)$$
[Wolf, Verstraete, Hastings, Cirac PRL 2008]

Entropy zero law or volume law = no correlations

What about a random state?

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The average entropy (random uniform) is close to maximal

[Page PRL 1995]





dimension of the subsystem

[Page PRL 1995, Bianchi, P.D. PRD 2019]

Sketch of the proof

step 1:

$$\langle \operatorname{Tr} \rho_A^r \rangle = \int d\mu(\psi) \operatorname{Tr} \rho_A^r = \int \left(\sum_{a=1}^{d_A} \lambda_a^r \right) \mu(\lambda_1, \dots, \lambda_{d_A}) \prod_{b=1}^{d_A} d\lambda_b$$

step 2:
 $\langle S_A \rangle = - \langle \operatorname{Tr} \rho_A \log \rho_A \rangle = -\lim_{r \to 1} \partial_r \langle \operatorname{Tr} \rho_A^r \rangle$
exact result:
 $\langle S_A \rangle = \Psi(d_A d_B + 1) - \Psi(d_B + 1) - \frac{d_A - 1}{2d_B}$
digamma function $\Psi = \Gamma'/\Gamma$
asymptotic:

$$\langle S_A \rangle \approx \log d_A - \frac{d_A^2 - 1}{2d_A d_B} \qquad d_B \gg 1$$

 $d_A \le d_B$

Sketch of the proof

[Bianchi, P.D. PRD 2019]

variance:

$$\Delta S_A^2 = \left\langle S_A^2 \right\rangle - \left\langle S_A \right\rangle^2$$

step 1:

$$\langle \mathrm{Tr} \rho_A^{r_1} \mathrm{Tr} \rho_A^{r_2} \rangle = \int \left(\sum_{a=1}^{d_A} \lambda_a^{r_1} \right) \left(\sum_{a=1}^{d_A} \lambda_a^{r_2} \right) \mu(\lambda_1, \dots, \lambda_{d_A}) \prod_{b=1}^{d_A} \mathrm{d}\lambda_b$$

step 2:

$$\left\langle S_A^2 \right\rangle = \lim_{\substack{r_1 \to 1 \\ r_2 \to 1}} \partial_{r_1} \partial_{r_2} \left\langle \mathrm{Tr} \rho_A^{r_1} \mathrm{Tr} \rho_A^{r_2} \right\rangle$$

exact result:

$$\Delta S_A^2 = \frac{d_A + d_B}{d_A d_B + 1} \Psi'(d_B + 1) - \Psi'(d_A d_B + 1) - \frac{(d_A - 1)(d_A + 2d_B - 1)}{4d_B^2(d_A d_B + 1)}$$

asymptotic:

$$\Delta S_A \approx \frac{1}{d_A d_B} \sqrt{\frac{d_A^2 - 1}{2}} \qquad d_B \gg 1 \qquad \qquad \Delta S_A \ll 1$$

The average entropy is typical!

 $d_A \le d_B$

 10^5 samples 200 bins $N_A = 4$ $N_B = 6$







Paramagnetic ionic salt



$$\label{eq:cuso_4} \begin{split} \underline{Cu}SO_4.K_2SO_4.6H_2O \\ \mu \simeq 0.9 \times 10^{-23} \; \mathrm{J/T} \\ \simeq 0.6 \times 10^{-4} \; \mathrm{eV/T} \end{split}$$



dimension of the subsystem

Typical entropy in presence of a center [Bianchi, P.D. PRD 2019]

Hilbert space structure:

$$\mathcal{H} = \bigoplus_{\zeta \in \mathcal{Z}} \left(\mathcal{H}_A^{(\zeta)} \bigotimes \mathcal{H}_B^{(\zeta)} \right)$$

$$|\psi\rangle = \sum_{\zeta} \sqrt{p_{\zeta}} \, |\phi_A^{(\zeta)}\rangle |\phi_B^{(\zeta)}\rangle$$

Entanglement entropy

$$S_A(\psi) = \sum_{\zeta} p_{\zeta} S_A(|\phi_A^{(\zeta)}\rangle |\phi_B^{(\zeta)}\rangle) - \sum_{\zeta} p_{\zeta} \log p_{\zeta}$$

The average entropy is typical

$$\langle S_A(\psi) \rangle = \sum_{\zeta} \frac{d_{A\zeta} d_{B\zeta}}{d} \langle S_{A\zeta} \rangle + \Psi(d+1) - \Psi(d_{A\zeta} d_{B\zeta} + 1)$$

Exact formula. Variance and other momenta also computed.



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$$\frac{1}{2} \left(\frac{\mathcal{C}(O_A, O_B)}{\|O_A\| \|O_B\|} \right)^2 \leq S_A(\psi) + \frac{S_B(\psi) - S_{AB}(\psi)}{\text{[Wolf, Verstraete, Hastings, Cirac PRL 2008]}}$$

random states (large dimension) follow a volume law = no correlations



Bell-network states

Gluing quantum polyhedra with entanglement [Bianchi, Baytas, Yokomizo, PRD 2018]

fluctuations of nearby quantum polyhedra are in general uncorrelated (twisted geometry)

Use squeezed vacua techniques to correlate shapes link by link [Bianchi, Hackl, Guglielmon, Yokomizo, PRD 2016]

Enlarge Hilbert space, squeeze the oscillators and project into \mathcal{H}_{Γ}



 $\mathcal{H}_{\Gamma} =$



Schwinger rep of SU(2)

$$\begin{split} \Gamma, \lambda_{\ell}, \mathcal{B} \rangle &= \sum_{j_{\ell}} \prod_{\ell} \left(1 - |\lambda_{\ell}|^2 \right) \lambda_{\ell}^{2j_{\ell}} \sqrt{2j_{\ell} + 1} \left| \Gamma, j_{\ell}, \mathcal{B} \right\rangle \quad \left| \Gamma, j_{\ell}, \mathcal{B} \right\rangle = \sum_{i_n} \mathcal{S}_{\Gamma}(j_{\ell}, i_n) \bigotimes_n |i_n\rangle \\ \lambda_{\ell} \in \mathbb{C} \quad \text{link squeezing parameter (average area, extrinsic curvature)} \quad \text{symbol of the graph} \end{split}$$

Bell-network states (analytic and numerics) $\mathcal{H}_{\Gamma} = \bigoplus_{j_{\ell}} \mathcal{L}_{\Gamma}$ $|\Gamma, j_{\ell}, \mathcal{B}\rangle = \sum_{i_n} \mathcal{S}_{\Gamma}(j_{\ell}, i_n) \bigotimes_n |i_n\rangle$ [Bianchi, P.D, Vilensky PRD 2019]

Asymptotic bound on the entanglement entropy

$$\left(\left|\partial A\right| - \frac{C_{\Gamma,A}}{2}\right)\log\lambda \le S_A \le \left(\left|\partial A\right| - \frac{3}{2}\right)\log\lambda$$

 $j_\ell \to \lambda j_\ell$

Numeric computations are available, at the moment, only for small graphs (but large spins)



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From numerical fit we can infer the bound on correlations

$$\frac{1}{2} \left(\frac{\mathcal{C}\left(O_A, O_B\right)}{\|O_A\| \|O_B\|} \right)^2 \le 0.06 \log \lambda$$

while for random states we obtain

$$\frac{1}{2} \left(\frac{\mathcal{C}(O_A, O_B)}{\|O_A\| \|O_B\|} \right)^2 \le O(\lambda^{-1})$$



 $j_\ell \to \lambda j_\ell$

Conclusions

Page curve and its variance:

- Random states have typical entropy
- Unlikely to have maximum entropy
- Concentration of measure
- Vanishing correlations
- Even in presence of a center
- Temperature from entanglement
- Half volume correction (heat capacity)

Entanglement entropy in a Bell-network state:

- Analytic asymptotics
- A rea-law from intertwiner entanglement
- Numerical code
- Non-vanishing intertwiner correlations

