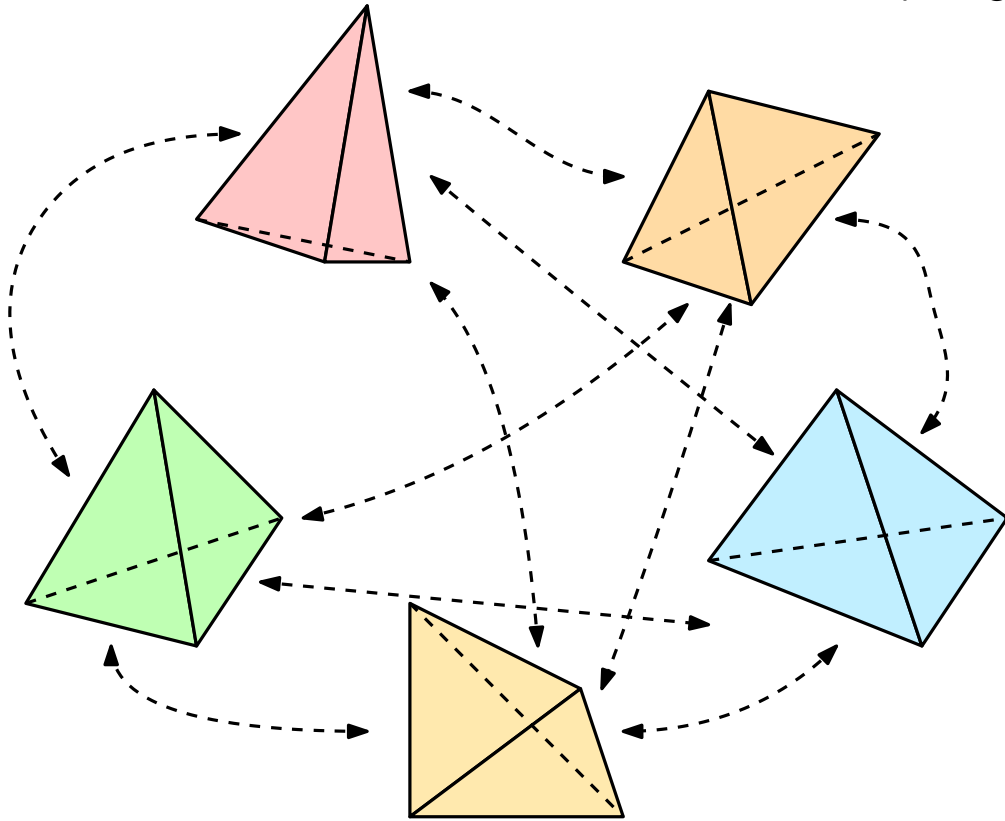


Pietro Donà

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The role of
local flatness
in spin foam
theories

EPRL spin foam theory



Dynamics for LQG as a path integral



Regularized on **simplicial** triangulations
(various extensions exists)



Connection with **discrete** GR



Numerical calculations are possible

EPRL spin foam theory



Dynamics for LQG as a path integral



Analytic calculations very complicated



Regularized on simplicial triangulations
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Where are Einstein equations?



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How?



Why?

Regge geometries emerge in the large quantum numbers of many spin foam models

Bivector reconstruction theorem

John Barrett's results (2010ish)



Amazing mathematical result. However...

- Not constructive (proof of existence)
- Mixes a lot of ingredients (hard to follow)
- Vertex amplitude specific (awkward extensions)
- Slight changes requires a complete rework (i.e. Muxin & co. 10 years ago)
- It is just **old**

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Different reasoning, similar conclusion, physically different models
(Topological BF $SU(2)$, EPRL Lorentzian and Euclidean, EPRL extensions, Barrett-Crane)

Geometry appears always with the same mechanism! Why?

Idea!



The emergence of Regge geometries does not depend on the details of the spin foam model. It comes from

Local flatness

An overlooked property...

A short math interlude

Spinors are not complicated!

They are an extremely powerful tool for calculations.

$$|z\rangle := \begin{pmatrix} z_0 \\ z_1 \end{pmatrix} \in \mathbb{C}^2$$

Using indices and tensors is more elegant, but I always mess up. I prefer Dirac's notation.

Complex structure

$$\langle z| := (\bar{z}_0, \bar{z}_1)$$

Scalar product

$$\langle w|z\rangle = \bar{w}_0 z_0 + \bar{w}_1 z_1$$

Dual spinor

$$|z] := \begin{pmatrix} -\bar{z}_1 \\ \bar{z}_0 \end{pmatrix}$$

A spinor and its dual $|z\rangle, |z]$ are linearly independent, orthogonal, have opposite chirality, they form a basis of the spinorial space, and have a very interesting geometrical interpretation

Geometrical interpretation

To simplify formulas I will work with norm 1 spinors $\langle z|z\rangle = 1$

Not a restriction. I am just lazy!

With each (unit) spinor we can build a (unit) vector

$$\langle z|\vec{\sigma}|z\rangle = -\vec{n}$$

independent from the phase of the spinor.



Geometrical interpretation

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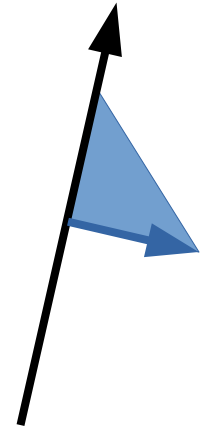
independent from the phase of the spinor.

There is more (flag/phase info)

(unit) frame vector

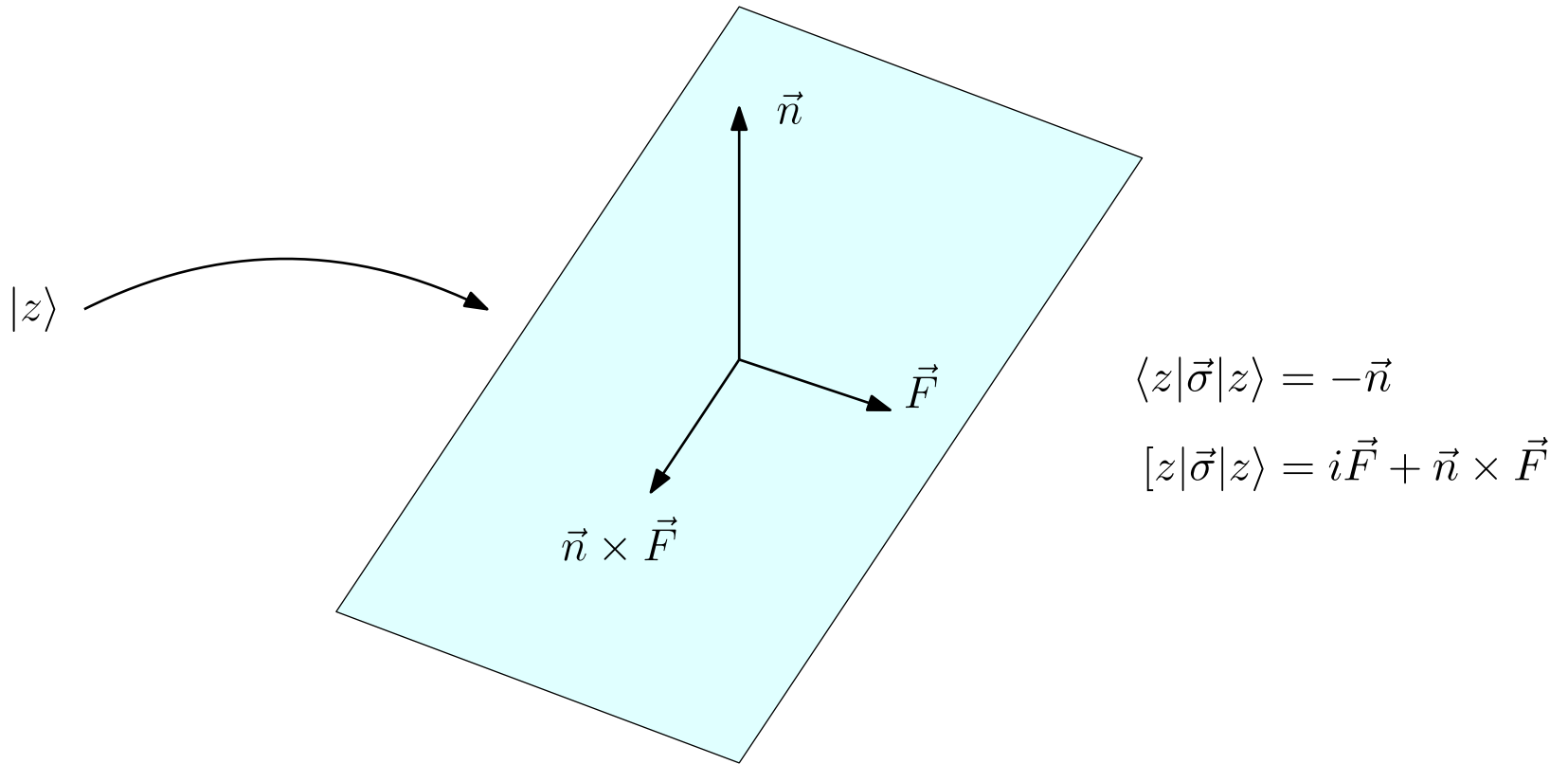
orthogonal to the vector

$$[z|\vec{\sigma}|z\rangle = i\vec{F} + \vec{n} \times \vec{F}$$



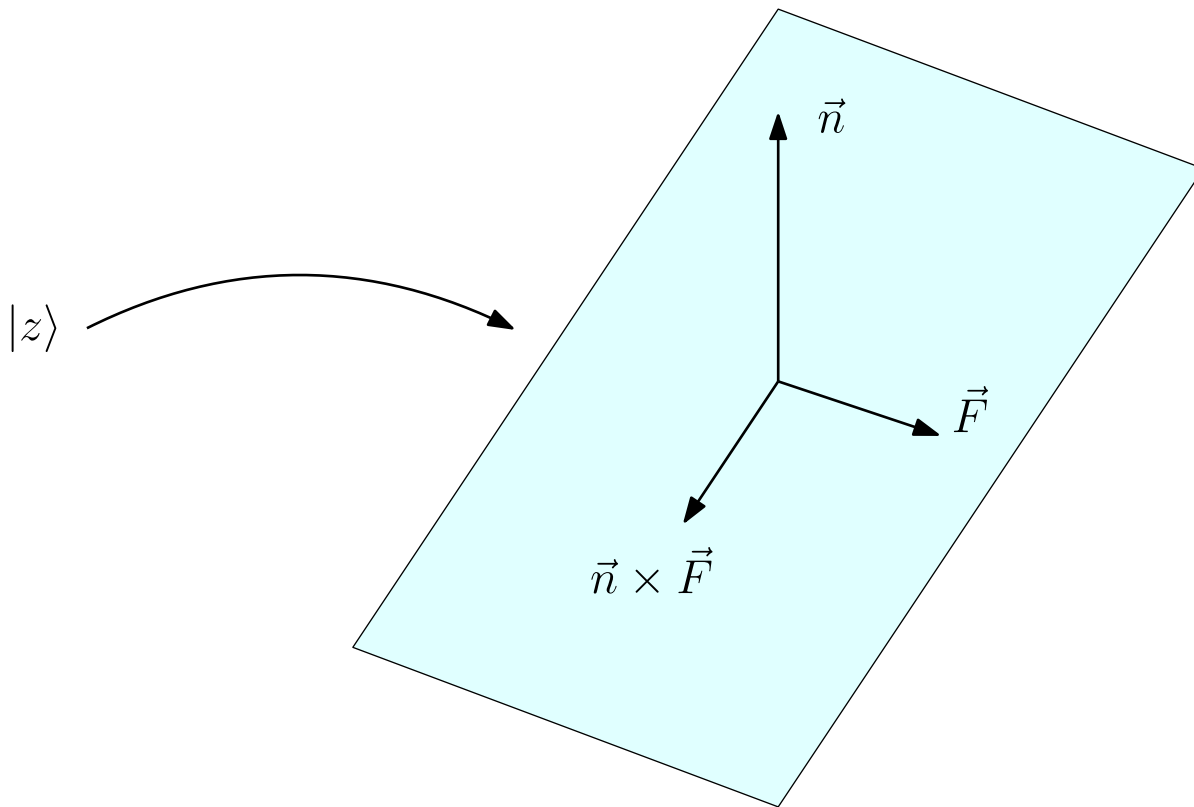
$\{\vec{n}, \vec{F}, \vec{n} \times \vec{F}\}$ orthonormal basis of Euclidean \mathbb{R}^3 (framed plane)

Spinors as framed planes



Framed plane in Euclidean 3D space

Spinors as framed planes



$$\langle z | \vec{\sigma} | z \rangle = -\vec{n}$$

$$[z | \vec{\sigma} | z] = i\vec{F} + \vec{n} \times \vec{F}$$

Déjà vu?

It is the same geometrical picture of LQG

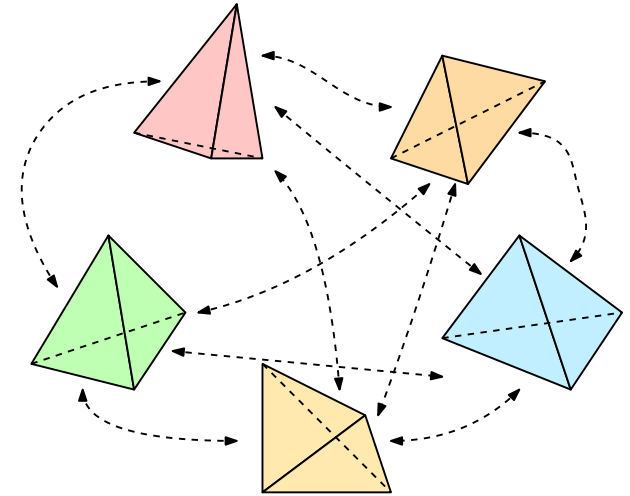
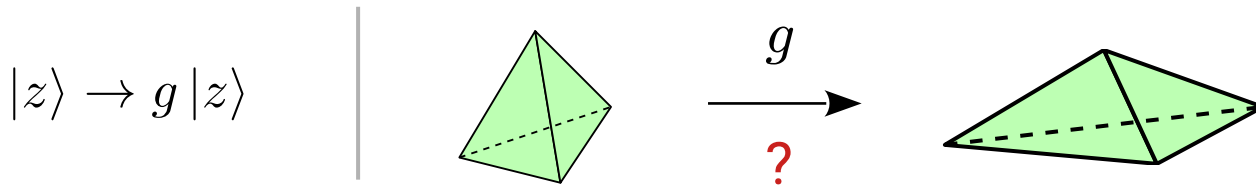
(twistorial phase space, twisted geometries - Simone, Etera, Laurent ...)



Why spinors as **framed planes**?

Direct interpretation in terms of LQG variables.
(twisted geometries are also parametrized in terms of spinors)

There is a **natural action** of Lorentz group on spinors!
(model independent, depends on the choice of **representation**, no extra ingredients)



(EPRL suggests electric part of gamma simple bivectors)

What can we infer from elementary properties?

Holonomies as map between **framed planes**

Clever parametrization of $SL(2, \mathbb{C})$

$$g = e^{\frac{\omega}{2}} |w\rangle\langle z| + e^{-\frac{\omega}{2}} |w][z|$$

A source spinor, a target spinor and a complex number (**redundant**, two extra phases, helps with interpretation)

Similar to twisted geometries parametrization of the Ashtekar holonomy. (particular adapted basis)

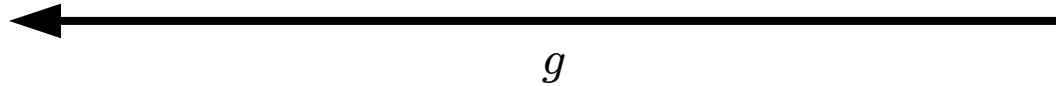
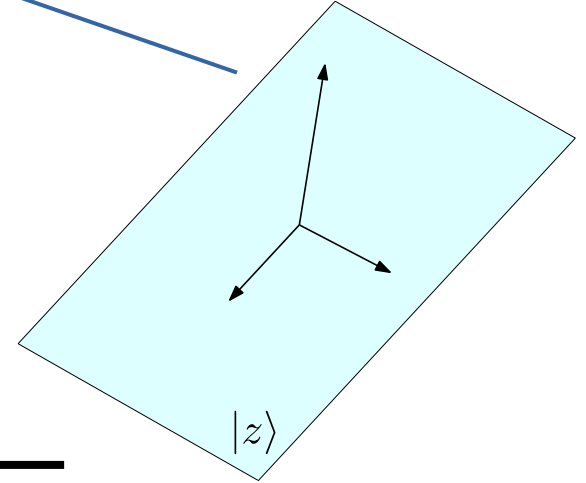
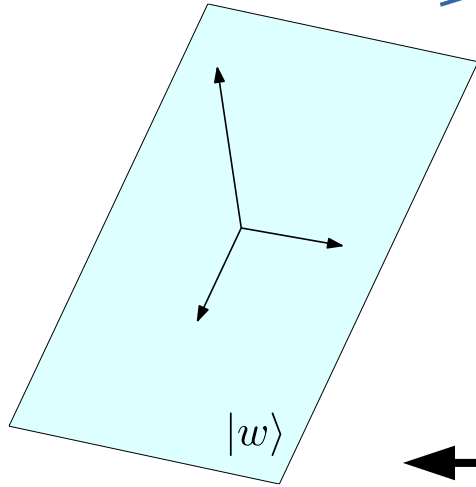
Geometry?

Holonomies as map between **framed planes**

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Target framed plane

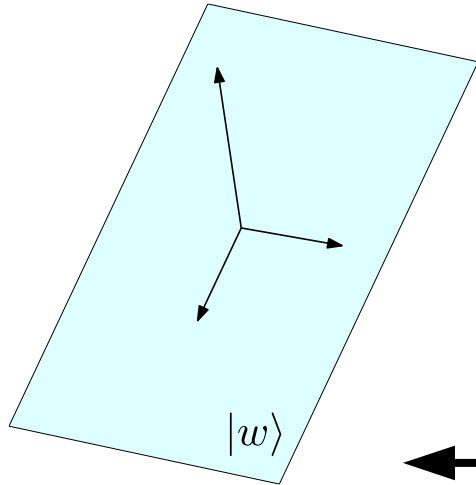
Source framed plane



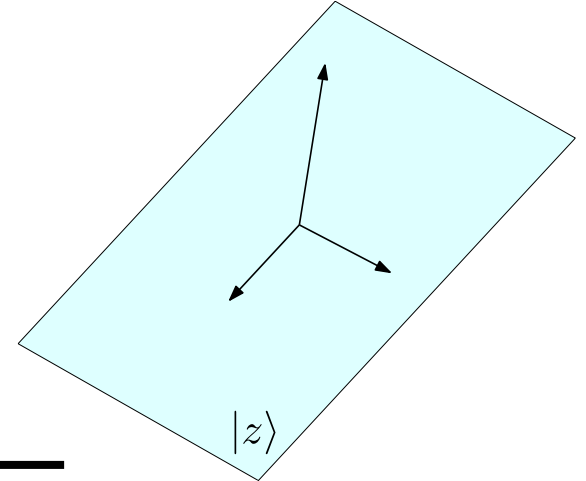
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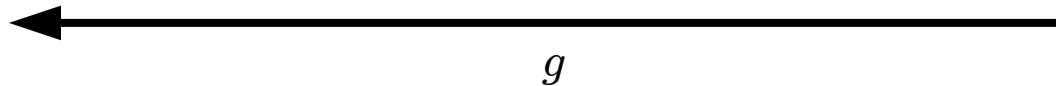


Source framed plane



Re ω Represent a boost between framed planes (**polar**)

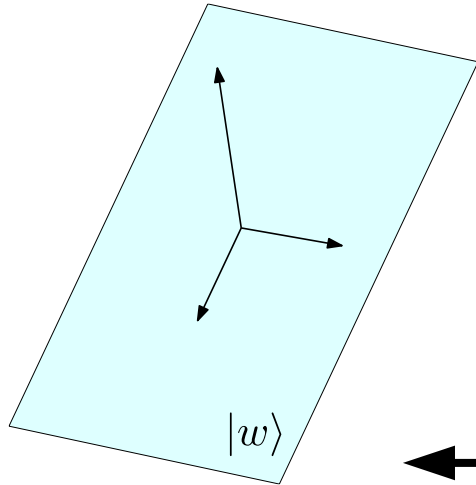
Im ω Represent a twist of the frames (**misalignment**)



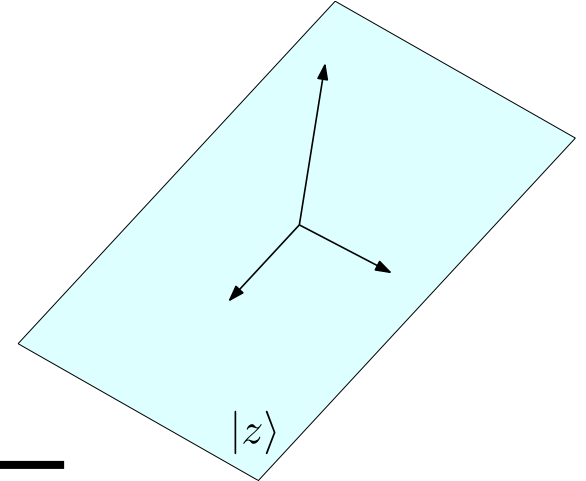
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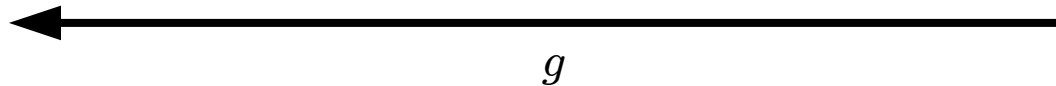


Source framed plane



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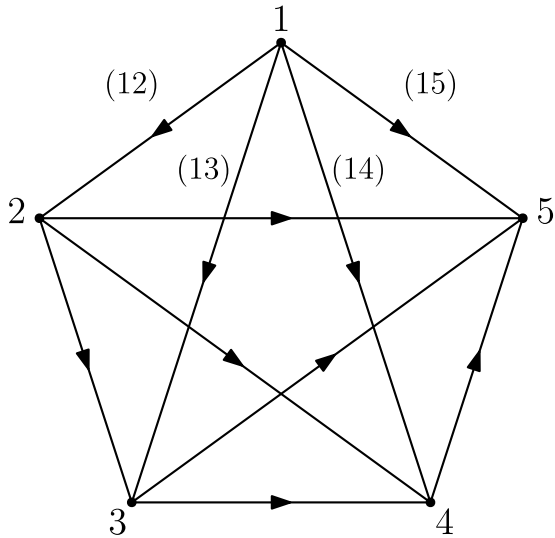


Geometry intrinsically associated to the holonomy

(similar to twisted geometries, but I am not talking about fluxes)

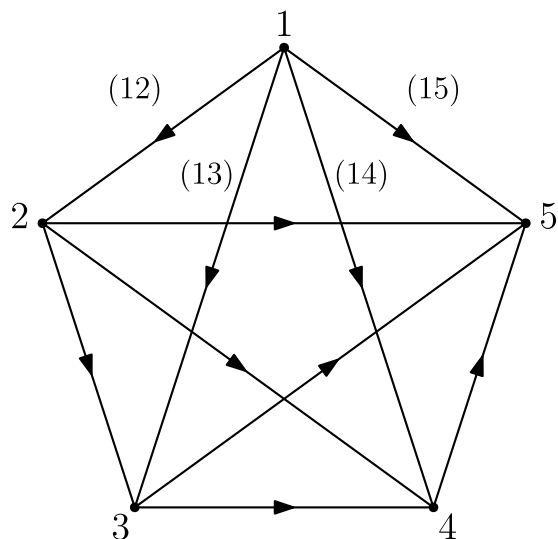
Setting the scene for **local flatness**

Consider the 2-complex of one 4-simplex (no geometric data, just **combinatorial** information)



Setting the scene for local flatness

Consider the 2-complex of one 4-simplex (no geometric data, just **combinatorial** information)



Five edges (dual to tetrahedra)

$a=1, 2, 3, 4, 5$

Ten faces (wedges) given by a couple of edges

$(ab)=(12), (13), (14), (15), \dots$

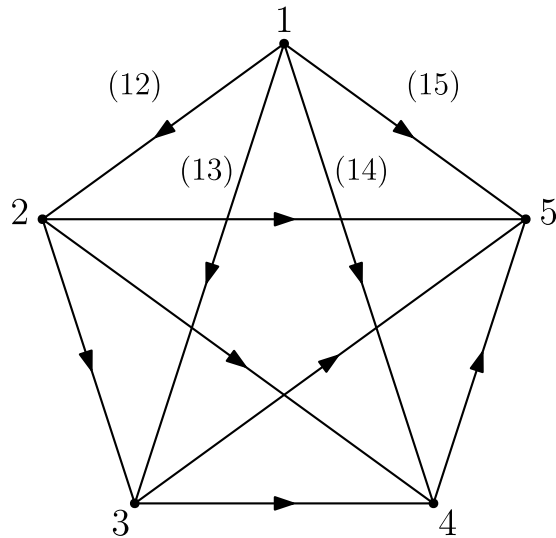
Associate one $SL(2, \mathbb{C})$ holonomy to each wedge

$$g_{ab} = e^{\frac{\omega_{ab}}{2}} |z_{ba}\rangle \langle z_{ab}| - e^{-\frac{\omega_{ab}}{2}} |z_{ba}\rangle [z_{ab}|$$

(**changed convention slightly**, helps with interpretation)

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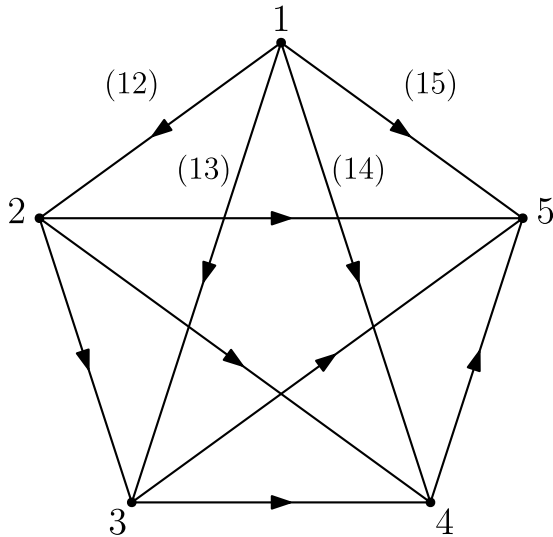
$$g_{ab} = e^{\frac{\omega_{ab}}{2}} |z_{ba}\rangle \langle z_{ab}| - e^{-\frac{\omega_{ab}}{2}} |z_{ba}\rangle [z_{ab}|$$

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Ten **holonomies** describe the **parallel transport** from one edge to another. They associate **four framed planes** to each edge and a **complex number** to each face.

Setting the scene for **local flatness**

Consider the 2-complex of one 4-simplex (no geometric data, just **combinatorial** information)



Tetrahedra?

Holonomies knows about angles.

No areas! No closure condition!

(SU(2) edge invariance. Info about fluxes. If confusing think about them as tetrahedra. They are one closure condition away)

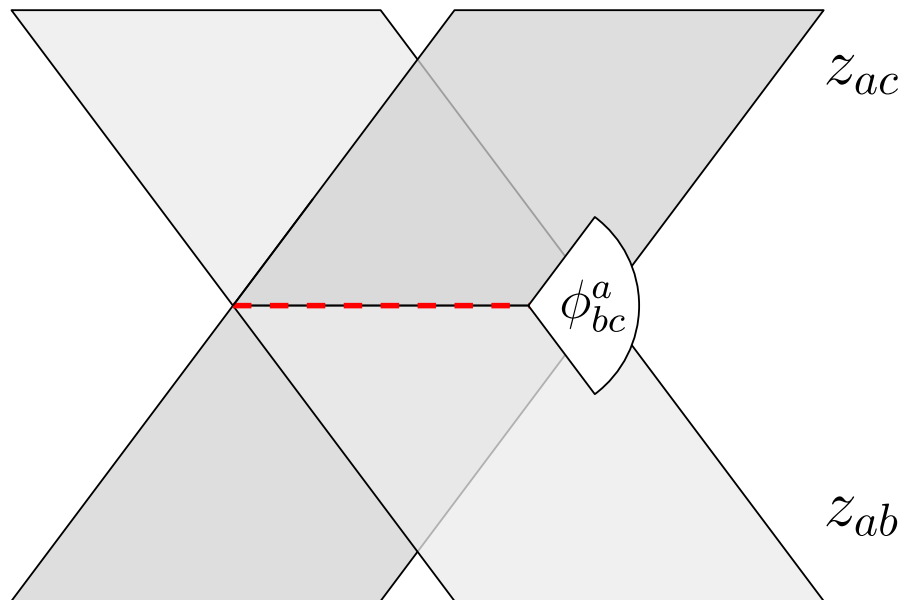
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Some geometric quantities

3d dihedral angles

(angle between 2 framed planes at the same edge)

$$\cos \phi_{bc}^a = \vec{n}_{ab} \cdot \vec{n}_{ac} = 2|\langle z_{ab} | z_{ac} \rangle|^2 - 1$$



Some geometric quantities

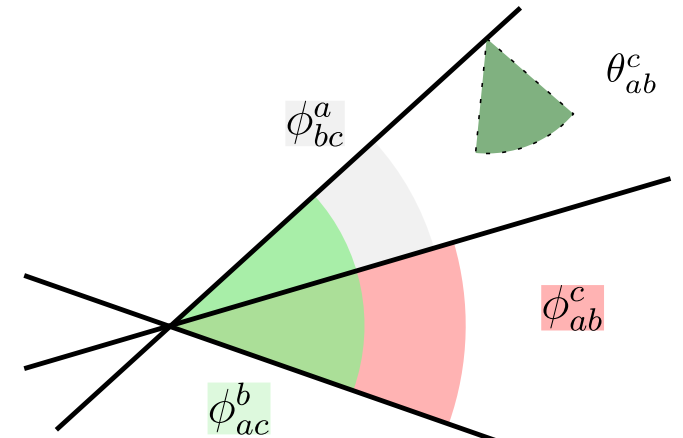
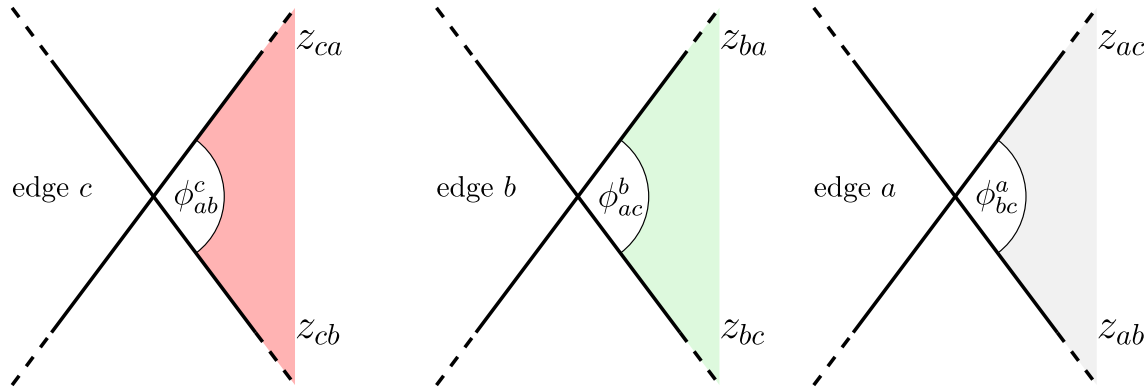
Spherical cosine law and sine law

(local **embedding** of 3D hyperplanes in 4D – signature?)

$$\cos \hat{\theta}_{ab}^c = \frac{-|\langle z_{ca} | z_{cb} \rangle|^2 + |\langle z_{ab} | z_{ac} \rangle|^2 |\langle z_{ba} | z_{bc} \rangle|^2 + |\langle z_{ac} | z_{ab} \rangle|^2 |\langle z_{ba} | z_{bc} \rangle|^2}{2|\langle z_{ac} | z_{ab} \rangle \langle z_{ac} | z_{ab} \rangle \langle z_{ba} | z_{bc} \rangle \langle z_{ba} | z_{bc} \rangle|} = \frac{\cos \phi_{ab}^c + \cos \phi_{bc}^a \cos \phi_{ac}^b}{\sin \phi_{bc}^a \sin \phi_{ac}^b}$$

$$\sin \phi_{ac}^b \sinh \hat{\theta}_{ab}^c = \sin \phi_{ab}^c \sinh \hat{\theta}_{ac}^b$$

(schematic picture in 1 lower dimension – I cannot draw in 4D)

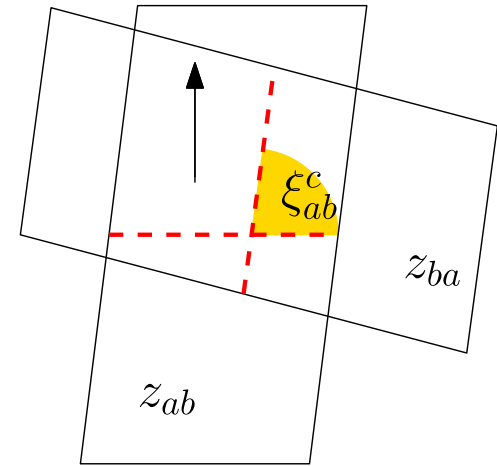
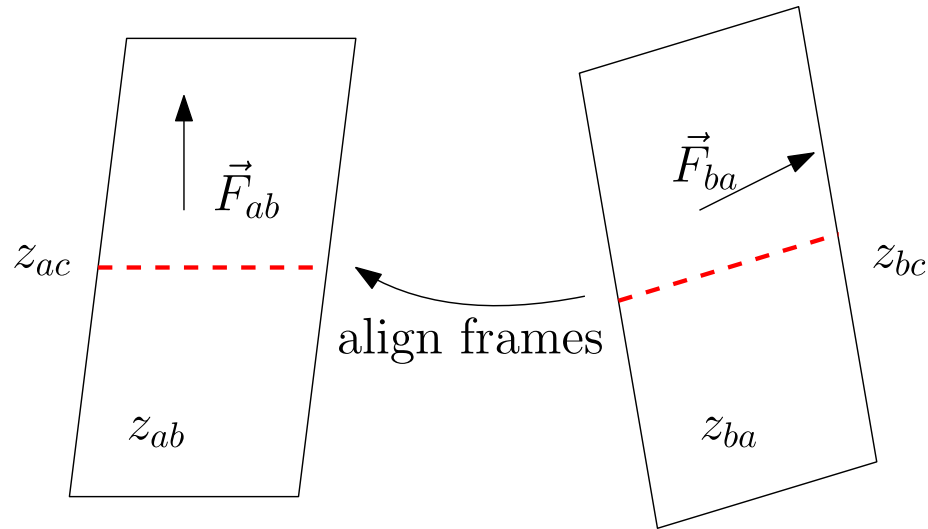


Some geometric quantities

Twist angle

(measure twist between frames using a third as reference -
the same one defined by Bianca and Jimmy or Simone and Fabio)

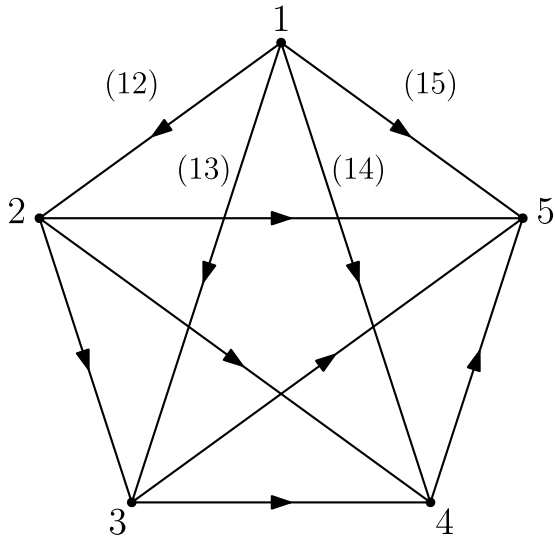
$$\xi_{ab}^c = \arg \left(\frac{\langle z_{ac} | z_{ab} \rangle \langle z_{ab} | z_{ac} \rangle}{\langle z_{bc} | z_{ba} \rangle \langle z_{ba} | z_{bc} \rangle} \right)$$



Holonomies and local flatness

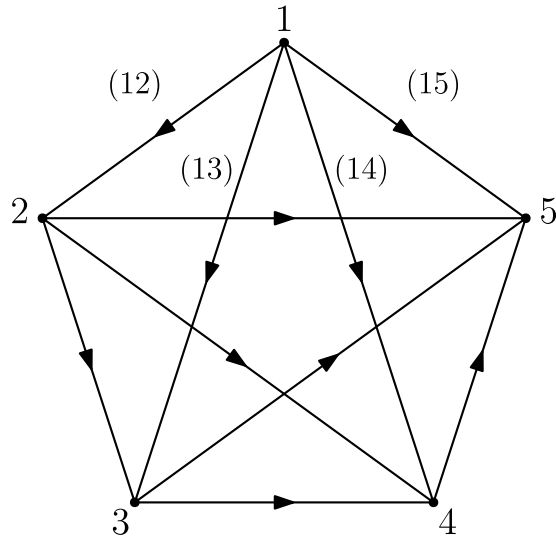
Flat building blocks!

We **require** each 4-simplex to be flat! The 2-complex is **locally flat**.



Holonomies and local flatness

Flat building blocks!



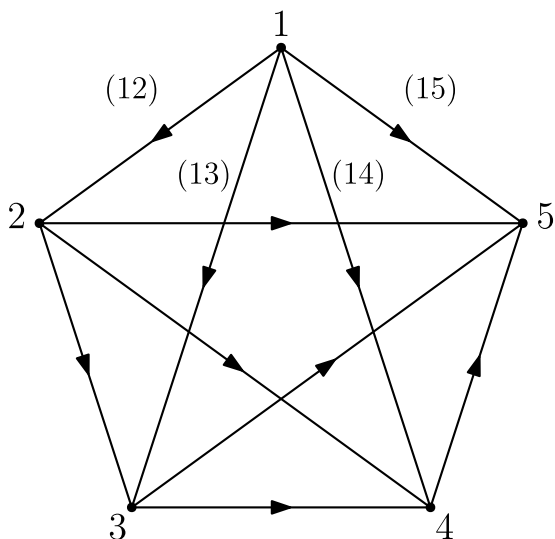
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$$g_{ca}g_{bc}g_{ab} = \mathbb{1}$$

Holonomies and local flatness

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In terms of holonomies the parallel transport on **every** closed **cycle** in the 4-simplex is trivial

$$g_{ca}g_{bc}g_{ab} = \mathbb{1}$$

Constraints on holonomies = constraints on the geometries (framed planes, spinors and complex angles)

How to solve? Smart projection on components

$$[z_{ac}|g_{ab}^{-1}|z_{bc}] = [z_{ac}|g_{ca}g_{bc}|z_{bc}]$$

$$\langle z_{ac}|g_{ab}^{-1}|z_{bc} \rangle = \langle z_{ac}|g_{ca}g_{bc}|z_{bc} \rangle$$

Plus other two. 4 complex scalar equations.

Combine to find...

Holonomies and local flatness

Local flatness

$$g_{ca}g_{bc}g_{ab} = \mathbb{1}$$

=

Complex angle determined by spinors

$$\cosh(\omega_{ab} + i\xi_{ab}^c) = \cos \hat{\theta}_{ab}^c$$

$$\sin \phi_{ac}^b \sinh(\omega_{ab} + i\xi_{ab}^c) = \sin \phi_{ab}^c \sinh(\omega_{ca} + i\xi_{ac}^b)$$

for every cycle = constrains also the spinors!

Solutions?

Studied by Me and Simone 2 years ago

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Solutions?

Studied by Me and Simone 2 years ago

Lorentzian sector

$$|\cos \hat{\theta}_{ab}^c| > 1$$

The other sector is the topological one (vector and euclidean) with SU(2) holonomies

Lorentzian geometries (need edge independence = angle matching = "shape" matching)

$$\omega_{ab} = \epsilon \theta_{ab} + i\epsilon \chi_{ab} \pi - i\xi_{ab}$$

orientation, dihedral angle, local causal structure, twist angles

Holonomies and local flatness

Local flatness

$$g_{ca}g_{bc}g_{ab} = \mathbb{1}$$

=

Complex angle determined by spinors

+

angle matching conditions (strongly)

(restriction to the Lorentzian sector)

$$\omega_{ab} = \epsilon\theta_{ab} + i\epsilon\chi_{ab}\pi - i\xi_{ab}$$

+

closure conditions
(areas up to a scale)

$$\sum_{b \neq a} j_{ab} \vec{n}_{ab} = 0$$

Information about fluxes, related to
edge SU(2) invariance

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Information about fluxes, related to
edge SU(2) invariance

In **ALL locally flat** Lorentzian spin foam models
(with edge SU(2) invariance)
shape matched Lorentzian 4-simplices emerge!

General! No amplitude, embedding map, semiclassical regime or critical point eqs.

Local flatness in the EPRL model

(all the spin foam models are locally flat)

$$A_v = \int \prod_a dg_a \delta(g_1) \prod_{ab} D_{j_{ab} \zeta_{ba} j_{ab} \zeta_{ab}}^{(\gamma j_{ab}, j_{ab})} (g_b^{-1} g_a)$$

(edge holonomies)

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(edge holonomies)

The EPRL model is locally flat (imposed strongly)

$$A_v = \int \left(\prod_{ab} dg_{ab} D_{j_{ab} \zeta_{ba} j_{ab} \zeta_{ab}}^{(\gamma j_{ab}, j_{ab})} \boxed{g_{ab}} \right) \mathcal{C}_{LF}(g_{ab}, \dots, g_{cd})$$

+

(wedge holonomies)

$$\mathcal{C}_{LF}(g_{ab}, \dots, g_{cd}) = \delta(g_{13}^{-1} g_{23} g_{12}) \delta(g_{14}^{-1} g_{24} g_{12}) \delta(g_{15}^{-1} g_{25} g_{12}) \cdot \\ \delta(g_{14}^{-1} g_{34} g_{13}) \delta(g_{15}^{-1} g_{35} g_{13}) \delta(g_{15}^{-1} g_{45} g_{14}) \cdot$$

What is the role of the critical point equations?

$$A_v = \int \left(\prod_{ab} dg_{ab} D_{j_{ab} \zeta_{ba} j_{ab} \zeta_{ab}}^{(\gamma j_{ab}, j_{ab})}(g_{ab}) \right) \mathcal{C}_{LF}(g_{ab}, \dots, g_{cd})$$

Coherent boundary data allow evaluation of the amplitude's integrals at the saddle point. \longrightarrow critical point equations

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Coherent boundary data allow evaluation of the amplitude's integrals at the saddle point. \longrightarrow critical point equations

Closure conditions of the boundary data \longrightarrow Framed planes at the same edge close forming a framed tetrahedron (areas given by the spins)
(consequence of the SU(2) invariance of the amplitude)

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the amplitude's integrals at the saddle point.

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of the boundary data
(consequence of the SU(2)
invariance of the amplitude)



Framed planes at the same edge close forming a framed tetrahedron (areas given by the spins)

Alignment equations



The spinors (framed planes) parametrizing the holonomies and the boundary ones are the same

What is the role of the critical point equations?

$$A_v = \int \left(\prod_{ab} dg_{ab} D_{j_{ab} \zeta_{ba} j_{ab} \zeta_{ab}}^{(\gamma j_{ab}, j_{ab})} (g_{ab}) \right) \mathcal{C}_{LF}(g_{ab}, \dots, g_{cd})$$

Saddle point \longrightarrow

Closure conditions
(boundary described as
framed tetrahedra)

+

Alignment equations
(holonomies spinors coincide
with boundary ones)

Action at the critical point

$$i\lambda \sum_{ab} j_{ab} (\gamma \text{Re } \omega_{ab} + \text{Im } \omega_{ab})$$

The connection with the Regge-Action happens
only on-shell of the **local flatness** conditions

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The connection with the Regge-Action happens
only on-shell of the **local flatness** conditions

focus to the Lorentzian sector

$$\omega_{ab} = \boxed{\epsilon \theta_{ab}} + \boxed{i \epsilon \chi_{ab} \pi} - \boxed{i \xi_{ab}}$$

4D dihedral angle between
the framed tetrahedra (ab)

Local causal
structure

Twist between the
framed tetrahedra (ab)

SKIP

Other stuffs

Many vertices?

Analysis on vertex
amplitudes independently
(local flatness + closure + alignment)

+

Extra alignment equations
(framed tetrahedra shared by different
vertices coincide!)

Summing over the spins

Constraining face
holonomies?

Singular support of the face distribution
(mystic result by Hellmann and Kaminski)

Naive flatness problem arises when
you combine
Local flatness + singular support +
alignment

Topological BF: $\delta(g_f) \longrightarrow g_f = \mathbb{1}$

EPRL: $f_{EPRL}(g_f) \longrightarrow g_f = e^{\frac{\omega_f}{2}} |\zeta\rangle\langle\zeta| + e^{-\frac{\omega_f}{2}} |\zeta][\zeta|$ with $\gamma \text{Re}\omega_f + \text{Im}\omega_f = 0 \pmod{4\pi}$

Conclusions

Local flatness is responsible of the local emergence of Regge geometry in spin foam models

Integrating over locally flat holonomies with $SU(2)$ edge invariance restricted on the Lorentzian sector = Summing over Lorentzian 4-simplices

(effective spin foam models and area-angle Regge calculus)

Quantum simplicity constraints (alignment + action)

Secondary simplicity constraints? (shape matching, imposed strongly)

Separation of ingredients is key to innovate (maybe new model? Simpler to do calculations! Top down construction! I have no concrete proposal.)